

Covariant Pencils of Matrices

AKA Wavelets

Vladimir V. Kisil

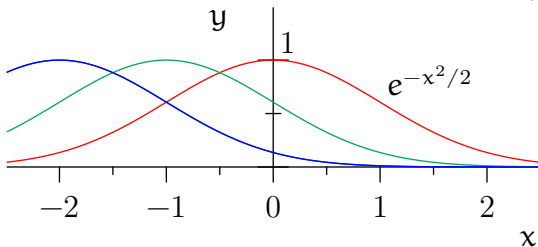
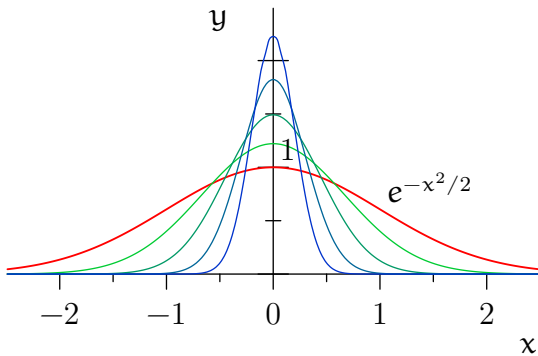
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Wavelets from the $ax + b$ group



Scaling and shift (the affine group) create a family of wavelets, which are used for wavelet transform:

$$\begin{aligned}\hat{f}(a, b) &= \langle f, \rho_{(a,b)} \phi_0 \rangle \\ &= \langle f, \phi_{(a,b)} \rangle\end{aligned}$$

The Gaussian $\phi(x) = e^{-x^2/2}$ as a mother wavelet produces an approximation of δ -function.

The mother wavelet $\phi(x) = \frac{1}{x+i}$ generates the Cauchy integral formula for the upper half-plane.

Wavelets and Groups

Let G be a group and ρ be its unitary irreducible representation in a Hilbert space H . For a fixed *mother wavelet* $\phi_0 \in H$ define *wavelet transform* from H to $C_b(G)$:

$$[\mathcal{W}f](a, b) = \langle f, \rho_{(a,b)} \phi_0 \rangle = \langle \rho_{(a,b)^{-1}} f, \phi_0 \rangle.$$

Let Λ be the *left regular representation* on G :

$$\Lambda(g) : f(h) \rightarrow f(g^{-1}h).$$

The following properties of the wavelet transform are of interest:

Proposition

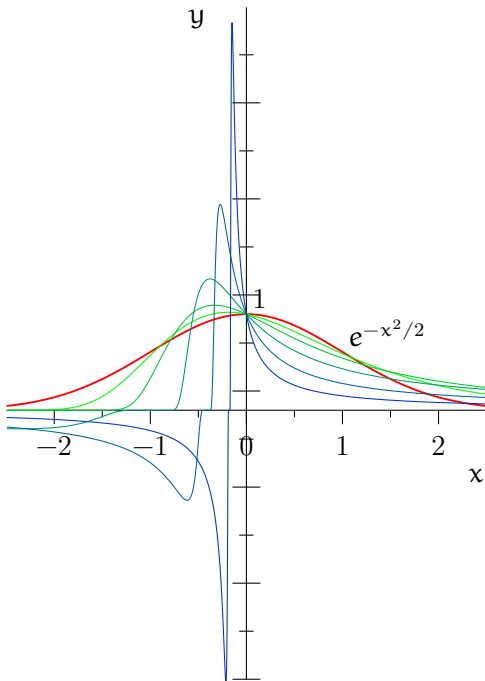
- 1 \mathcal{W} intertwines ρ and Λ :

$$\mathcal{W}\rho(g) = \Lambda(g)\mathcal{W}.$$

- 2 The image of \mathcal{W} is invariant under left shifts.

- 3 The image of \mathcal{W} is generated by images of mother wavelets $\mathcal{W}\phi_0$.





If we extend the group from $ax + b$ to $SL_2(\mathbb{R})$ then the same Gaussian as the mother wavelet will produce not only an approximation of the δ -function but also an approximation of the δ' distribution as well.

However, the extension of the group will not affect the mother wavelet $\frac{1}{x+i}$; it will still generate the Cauchy integral, because it is an eigenvector of the subgroup K .

Thus, choices of group and mother wavelets produce different frameworks.

The group $SL_2(\mathbb{R})$ and operator valued representations

The group $G = SL_2(\mathbb{R})$ consists of 2×2 matrices $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$ with real entries and the unit determinant $ad - bc = 1$.

Let H be a Hilbert space, for a bounded linear operator T on H we define a paraphrase of the resolvent:

$$R(g) = (cT + dI)^{-1} \quad \text{where} \quad g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{R}).$$

Obviously $R(g)$ contains the essential information about the operator T .
How to extract it?

Consider the space $L(G)$ of functions spanned by all left translations of $R(g)$. The left action $g : f(h) \mapsto f(g^{-1}h)$ of $SL_2(\mathbb{R})$ on this space is a linear representation of this group.

Vector and scalar versions

Function $R(g)$ contain too much information, we may restrict it to get a more detailed view. For vectors $u, v \in H$ we also consider vector and scalar-valued functions related to resolvent:

$$R_u(g) = (cT + dI)^{-1}u, \quad \text{and} \quad R_{u,v}(g) = \langle (cT + dI)^{-1}u, v \rangle$$

where $(cT + dI)^{-1}u$ is understood as a solution w of the equation $u = (cT + dI)w$ if it is unique.

Example

$$T_- = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad T_0 = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \quad T_+ = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$
$$R_-(g) = \begin{pmatrix} d & -c \\ c & d \end{pmatrix} \quad R_0(g) = \begin{pmatrix} d & -c \\ -c & d \end{pmatrix} \quad R_+(g) = \begin{pmatrix} d & -c \\ 0 & d \end{pmatrix}$$

Note, $T_\sigma = \sigma I$, that is T_σ is a model for hypercomplex imaginary unit.

Covariant Calculus

Analytic function theory in the upper half-plane \mathbb{R}_+^2 is mainly a theory of the *discrete series* representation of $SL_2(\mathbb{R})$ group of 2×2 matrices:

$$\rho_m(g) : f(z) \mapsto \frac{1}{(cz + d)^m} f\left(\frac{az + b}{cz + d}\right), \quad g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{R}). \quad (1)$$

To get a definition of functional calculus we replace of a *homomorphism property* by a *symmetric covariance*. One possible realisation is:

Definition

An analytic functional calculus for an element $T \in \mathfrak{A}$ and an \mathfrak{A} -module M is a *continuous linear* mapping $\Phi : A(\mathbb{R}_+^2) \rightarrow A(\mathbb{R}_+^2, M)$ such that Φ is an *intertwining operator*

$$\Phi \rho_1 = \rho_T \Phi$$

between two representations of the $SL_2(\mathbb{R})$ group ρ_1 (1) and ρ_T , which is the restriction of the right shifts to the space generated by $R_v(g)$.

Covariant Spectrum

as the Support of the Calculus

Definition

A covariant spectrum is defined as the support of covariant functional calculus, that is a decomposition of the intertwining map into primary subrepresentations.

Such components associated with the above $SL_2(\mathbb{R})$ calculus fall into three large classes, with the pairings defined by the usual and indefinite inner products as well as degenerate one.

Example

For the above matrices T_σ the irreducible components are isomorphic to analytic spaces of hypercomplex functions under the fraction-linear transformations:

$$w \mapsto \frac{aw + b}{cw + d}, \quad \text{where } w = u + iv, \text{ and } v^2 = -1, 0, 1.$$

$SL_2(\mathbb{R})$ and Its Subgroups

$SL_2(\mathbb{R})$ is the group of 2×2 matrices with real entries and $\det = 1$. A two dimensional subgroup F (F') of lower(upper) triangular matrices:

$$F = \left\{ \frac{1}{\sqrt{a}} \begin{pmatrix} a & 0 \\ c & 1 \end{pmatrix} \right\}, \quad F' = \left\{ \frac{1}{\sqrt{a}} \begin{pmatrix} a & b \\ 0 & 1 \end{pmatrix} \right\}, \quad a \in \mathbb{R}_+, b, c \in \mathbb{R}.$$

F is the group of affine transformations of the real line ($ax + b$ group), isomorphic to the upper half-plane.

There are also three one dimensional continuous subgroups:

$$A = \left\{ \begin{pmatrix} e^t & 0 \\ 0 & e^{-t} \end{pmatrix} = \exp \begin{pmatrix} t & 0 \\ 0 & -t \end{pmatrix}, t \in \mathbb{R} \right\}, \quad (2)$$

$$N = \left\{ \begin{pmatrix} 1 & t \\ 0 & 1 \end{pmatrix} = \exp \begin{pmatrix} 0 & t \\ 0 & 0 \end{pmatrix}, t \in \mathbb{R} \right\}, \quad (3)$$

$$K = \left\{ \begin{pmatrix} \cos t & \sin t \\ -\sin t & \cos t \end{pmatrix} = \exp \begin{pmatrix} 0 & t \\ -t & 0 \end{pmatrix}, t \in (-\pi, \pi] \right\}. \quad (4)$$

... and Nothing Else

(up to a conjugacy)

Proposition

Any one-parameter subgroup of $SL_2(\mathbb{R})$ is conjugate to either A , N or K .

Proof.

Any one-parameter subgroup is obtained through the exponentiation

$$e^{tX} = \sum_{n=0}^{\infty} \frac{t^n}{n!} X^n \quad (5)$$

of an element X of the Lie algebra \mathfrak{sl}_2 of $SL_2(\mathbb{R})$. Such X is a 2×2 matrix with the zero trace. The behaviour of the Taylor expansion (5) depends from properties of powers X^n . This can be classified by a straightforward calculation. □

Elliptic, Parabolic, Hyperbolic

the First Appearance

Lemma

The square X^2 of a traceless matrix $X = \begin{pmatrix} a & b \\ c & -a \end{pmatrix}$ is the identity matrix times

$a^2 + bc = -\det X$. The factor can be negative, zero or positive, which corresponds to the three different types of the Taylor expansion (5) of e^{tX} .

It is a simple exercise in the Gauss elimination to see that through the matrix similarity we can obtain from X a generator

- of the subgroup K if $(-\det X) < 0$;
- of the subgroup N if $(-\det X) = 0$;
- of the subgroup A if $(-\det X) > 0$.

The determinant is invariant under the similarity, thus these cases are distinct.

$SL_2(\mathbb{R})$ and Homogeneous Spaces

Let G be a group and H be its closed subgroup.

The *homogeneous space* G/H from the equivalence relation: $g' \sim g$ iff $g' = gh, h \in H$. The *natural projection* $p : G \rightarrow G/H$ puts $g \in G$ into its equivalence class.

A continuous section $s : G/H \rightarrow G$ is a right inverse of p , i.e. $p \circ s$ is an identity map on G/H . Then the *left action* of G on itself:

$$\Lambda(g) : g' \mapsto g^{-1} * g', \quad \text{generates} \quad \begin{array}{ccc} G & \xrightarrow{g^*} & G \\ s \updownarrow p & & s \updownarrow p \\ G/H & \xrightarrow{g \cdot} & G/H \end{array}$$

If $G = SL_2(\mathbb{R})$ and $H = F$, then $SL_2(\mathbb{R})/F \sim \mathbb{R}$ and $p : \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mapsto \frac{b}{d}$:

$$s : u \mapsto \begin{pmatrix} 1 & u \\ 0 & 1 \end{pmatrix}, \quad g : u \mapsto p(g^{-1} * s(u)) = \frac{au + b}{cu + d}, \quad g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

$SL_2(\mathbb{R})$ as a Source of Imaginary Units

Consider $G = SL_2(\mathbb{R})$ and H be any subgroup in the Iwasawa decomp:

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} \alpha & 0 \\ 0 & \alpha^{-1} \end{pmatrix} \begin{pmatrix} 1 & v \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix}. \quad (6)$$

A right inverse s to the natural projection $p : G \rightarrow G/H$:

$$s : (u, v) \mapsto \frac{1}{\sqrt{v}} \begin{pmatrix} v & u \\ 0 & 1 \end{pmatrix}, \quad (u, v) \in \mathbb{R}^2, \text{ in the diagram}$$

$$\begin{array}{ccc} G & \xrightarrow{g^*} & G \\ s \updownarrow p & & s \updownarrow p \\ G/H & \xrightarrow{g_*} & G/H \end{array}$$

defines the G -action $g \cdot x = p(g \cdot s(x))$ on the homogeneous space G/H :

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} : (u, v) \mapsto \left(\frac{(au + b)(cu + d) - \sigma cv^2}{(cu + d)^2 - \sigma (cv)^2}, \frac{v}{(cu + d)^2 - \sigma (cv)^2} \right).$$

This becomes a Möbius map in (hyper)complex numbers:

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} : w \mapsto \frac{aw + b}{cw + d}, \quad w = u + iv, \quad i^2 (:= \sigma) = -1, 0, 1.$$

Möbius Transformations of \mathbb{R}^2

For *all* numbers define Möbius' transformation of $\mathbb{R}^2 \rightarrow \mathbb{R}^2$,
(in elliptic and parabolic cases this is even $\mathbb{R}_+^2 \rightarrow \mathbb{R}_+^2$):

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} : u + iv \mapsto \frac{a(u + iv) + b}{c(u + iv) + d}. \quad (7)$$

Product $\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} \tau & 0 \\ 0 & \tau^{-1} \end{pmatrix} \begin{pmatrix} 1 & \nu \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix}$ gives *Iwasawa*
 $SL_2(\mathbb{R}) = \mathbf{A}\mathbf{N}\mathbf{K}$. In all \mathbf{A} subgroups \mathbf{A} and \mathbf{N} acts uniformly:

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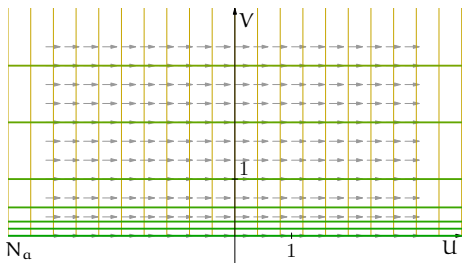
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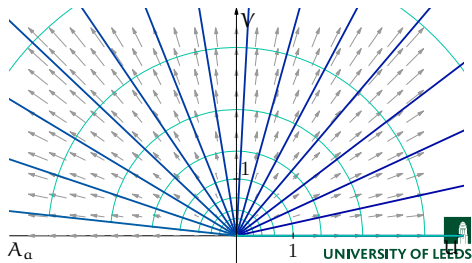
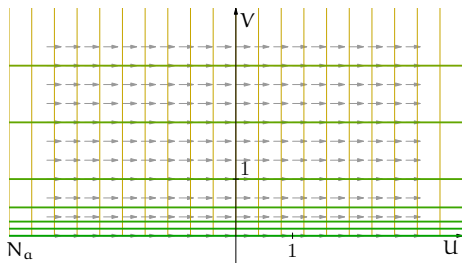


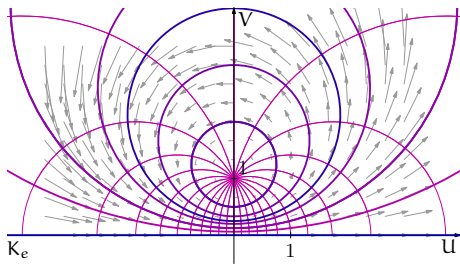
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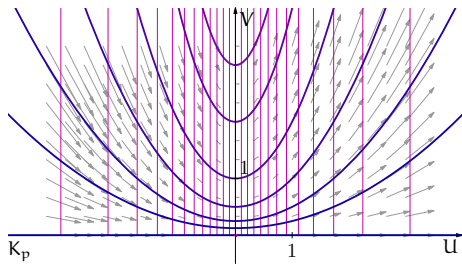
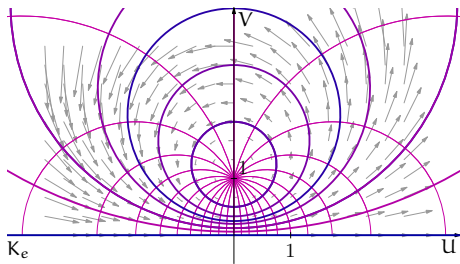
$$dK_e(u, v) = (1 + u^2 - v^2, \quad 2uv)$$

$$dK_p(u, v) = (1 + u^2, \quad 2uv)$$

$$dK_h(u, v) = (1 + u^2 + v^2, \quad 2uv)$$

$$dK_\sigma(u, v) = (1 + u^2 + \sigma v^2, \quad 2uv)$$

Figure: Depending from $i^2 = \sigma$ the orbits of subgroup K are circles, parabolas and hyperbolas passing $(0, t)$ with the equation $(u^2 - \sigma v^2) + v(\sigma t - t^{-1}) + 1 = 0$. This leads to elliptic, parabolic and hyperbolic analytic functions from induced representations.



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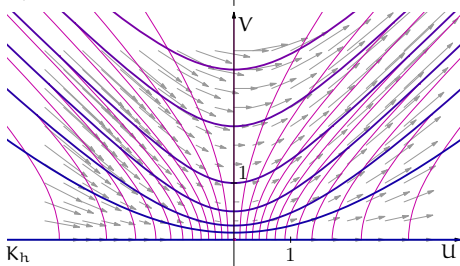
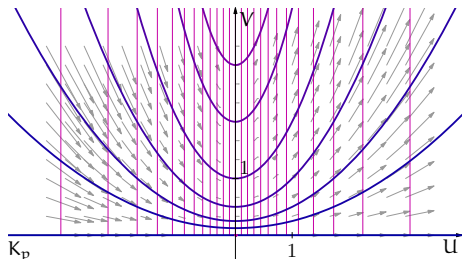
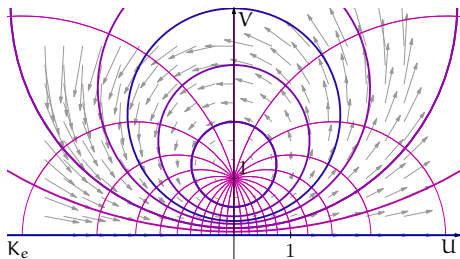
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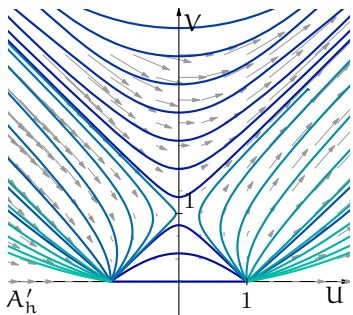
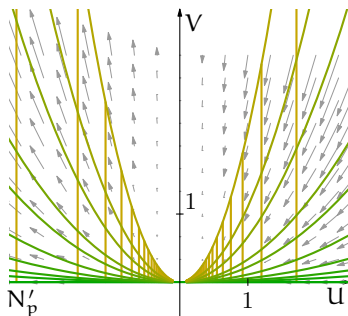
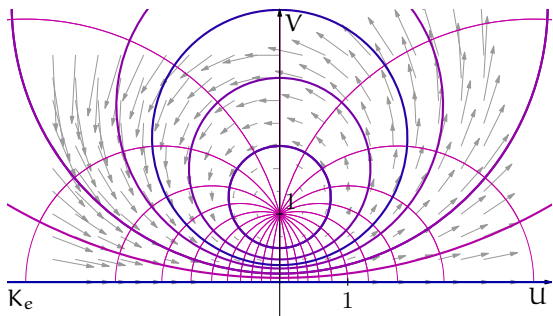
Elliptic, Parabolic, Hyperbolic

Möbius transformations with the three number systems produce three types of non-Euclidean geometries in the spirit of the Erlangen programme of F. Klein:

Numbers	complex	dual	double
Imaginary unit	$i^2 = -1$	$\epsilon^2 = 0$	$\epsilon^2 = 1$
Inv. objects	circles	parabolas	hyperbolas
$ z ^2 = z\bar{z}$ is	$x^2 + y^2$	x^2	$x^2 - y^2$
Inv. distance	$\sinh^{-1} \frac{ z-w }{2\sqrt{\Im[z]\Im[w]}}$	$\frac{ z-w }{2\sqrt{\Im[z]\Im[w]}}$	$\sin^{-1} \frac{ z-w }{2\sqrt{\Im[z]\Im[w]}}$
Equidist.	circles	parabolas	hyperbolas
Geodesics	circles	parabolas	hyperbolas

There are many beautiful geometrical results going through all three cases which we will skip in this talk.

Question: can this parallel be extended to analytic functions and operator theory?



Fix subgroups of $(0, 1)$ are:

$$K = \begin{pmatrix} \cos t & \sin t \\ -\sin t & \cos t \end{pmatrix} = \exp t \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} -$$

the elliptic case;

$$N' = \begin{pmatrix} 1 & 0 \\ t & 1 \end{pmatrix} = \exp t \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} - \text{the}$$

parabolic case;

$$A' = \begin{pmatrix} \cosh t & \sinh t \\ \sinh t & \cosh t \end{pmatrix} = \exp t \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} -$$

the hyperbolic case.

Induced Representations

Let G be a group, H its closed subgroup, χ be a linear representation of H in a space V . The set of V -valued functions with the property

$$F(gh) = \chi(h)F(g),$$

is invariant under left shifts.

The restriction of the left regular representation to this space is called an *induced representation*.

Equivalently we consider the *lifting* of $f(x)$, $x \in X = G/H$ to $F(g)$:

$$F(g) = \chi(h)f(p(g)), \quad p: G \rightarrow X, \quad g = s(x)h, \quad p(s(x)) = x.$$

This is a 1-1 map which transform the left regular representation on G to the following action:

$$[\rho'(g)f](x) = \chi(h)f(g \cdot x), \quad \text{where } gs(x) = p(gs(x))h.$$

In the case of $SL_2(\mathbb{R})$ we have three different types of actions.

Affine Group

For $G = \mathrm{SL}_2(\mathbb{R})$ and $H = \mathbb{F}$ the action on G/H is:

$$g : u \mapsto p(g^{-1} * s(u)) = \frac{au + b}{cu + d}, \quad \text{where } g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

We calculate also that

$$r(g^{-1} * s(u)) = \begin{pmatrix} (cu + d)^{-1} & 0 \\ c & cu + d \end{pmatrix}.$$

A generic character of \mathbb{F} is a power of its diagonal element:

$$\rho_\kappa \left(\begin{pmatrix} a & 0 \\ c & a^{-1} \end{pmatrix} \right) = a^\kappa.$$

Thus the corresponding realisation of induced representation is:

$$\rho_\kappa(g) : f(u) \mapsto \frac{1}{(cu + d)^\kappa} f \left(\frac{au + b}{cu + d} \right) \quad \text{where } g^{-1} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

Induced Wavelet Transform

Let $v_0 \in \mathcal{H}$ be an eigenfunction as follows:

$$\rho(\mathbf{h})v_0 = \chi(\mathbf{h}) \cdot v_0, \quad \text{for all } \mathbf{h} \in \tilde{\mathbf{H}}.$$

It is suitable to be the *mother wavelet* (*vacuum vector*). Then we have

$$\begin{aligned} [\mathcal{W}f](g\mathbf{h}) &= \langle f, \rho(g\mathbf{h})v_0 \rangle = \langle f, \rho(g)\rho(\mathbf{h})v_0 \rangle \\ &= \langle f, \chi(\mathbf{h}) \cdot \rho(g)v_0 \rangle = \chi(\mathbf{h}) \langle f, \rho(g)v_0 \rangle. \end{aligned}$$

For v_0 the *induced wavelet transform* $\mathcal{W} : \mathcal{H} \rightarrow L_\infty(G/H)$ by

$$[\mathcal{W}f](w) = \langle f, \rho_0(s(w))v_0 \rangle, \tag{8}$$

where $w \in G/\tilde{\mathbf{H}}$ and $s : G/\tilde{\mathbf{H}} \rightarrow G$.

It intertwines ρ with a representation induced by χ of $\tilde{\mathbf{H}}$.

Cauchy Integral Formula

Eigenvector of K

The infinitesimal version of the eigenvector property $\rho(\mathbf{h})\mathbf{v}_0 = \chi(\mathbf{h}) \cdot \mathbf{v}_0$ is $d\rho_K \mathbf{v}_0 = \lambda \mathbf{v}_0$, explicitly

$$nuf(u) + f'(u)(1 + u^2) = \lambda f(u).$$

The generic solution is:

$$f(u) = \frac{1}{(1 + u^2)^{n/2}} \left(\frac{u + i}{u - i} \right)^{i\lambda/2} = \frac{(u + i)^{(i\lambda - n)/2}}{(u - i)^{(i\lambda + n)/2}}.$$

To avoid multivalent function we need to put $\lambda = im$ and the minimal weight condition suggests $m = n$. The induced wavelet transform is:

$$\hat{f}(x, y) = \langle f, \rho_n f_0 \rangle = \int_{\mathbb{R}} f(u) \frac{\sqrt{y}}{u - x - iy} dx = \sqrt{y} \int_{\mathbb{R}} f(u) \frac{dx}{u - (x + iy)}$$



Other Integral Transforms

Eigenvalues of N

The subgroup N consists of shifts, the eigenfunction is $e^{\lambda u}$ and the induced wavelet transform coincides with the Fourier transform.

For the subgroup N' the derived representation is

$$d\rho^{N'} = (un) \cdot I - u^2 \cdot \partial_u.$$

The corresponding eigenvector is $f_0 = u^n e^{\frac{\lambda}{u}}$.

The induced representation

$$\hat{f}(x, y) = \langle f, \rho_n f_0 \rangle = \int_{\mathbb{R}} f(u) \frac{(u - x - iy)^n}{\sqrt{y}^n} e^{\frac{\lambda\sqrt{y}}{u-x-iy}} dx.$$

(It is a Fourier transform combined with the inversion).

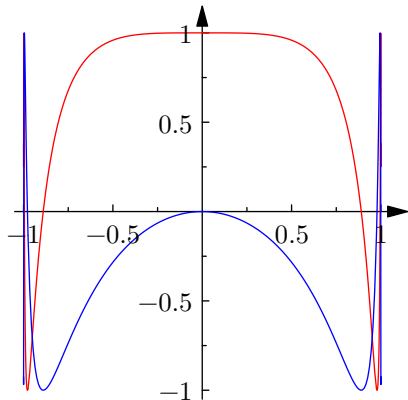
Other Integral Transforms

Eigenvalues of A

For the subgroup A the derived representation $d\rho_f^{A'}(u) = -nuf(u) + (u^2 - 1)f'(u)$. It has two singular point ± 1 .

$$\begin{aligned} f(x) &= \frac{1}{(u^2 - 1)^{n/2}} \left(\frac{u+1}{u-1} \right)^{\lambda/2} \\ &= \frac{(u+1)^{(\lambda-n)/2}}{(u-1)^{(\lambda+n)/2}}. \end{aligned}$$

The solution has compact support. Its meaning and applicability shall be investigated.



Bergman Integrals

as Wavelet Transforms in the Half Plane

Considering: $\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \frac{1}{\sqrt{v}} \begin{pmatrix} v & u \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{pmatrix}$ we chose:

- Eigenvector of K vacuum vector $v_0(z) = \frac{1}{(x+i)^m}$.
- wavelets or coherent states $v_m(g, z) = \rho_m(g)v_0(z) = (x - (u + iv))^{-m}$. They depend only from $u + iv \in \mathbb{R}_+^2$.
- The universally defined wavelet transforms

$$\mathcal{W}_m : f(z) \mapsto \mathcal{W}_m f(u) = \langle f(z), \rho_m v_0(u, z) \rangle$$

Then $v_m(u, z)$ are the Cauchy and Bergman kernels. The corresponding wavelet transform became Cauchy and Bergman integrals:

$$\mathcal{W}_1 f(u) = \frac{1}{2\pi i} \int_{\mathbb{R}} f(z) \frac{1}{u-z} dz,$$

$$\mathcal{W}_m f(u) = \int_{\mathbb{R}_+^2} f(z) \frac{1}{(\bar{u}-z)^m} (\mathcal{J}z)^{m-2} dz.$$

Fix Subgroups of i and ϵ

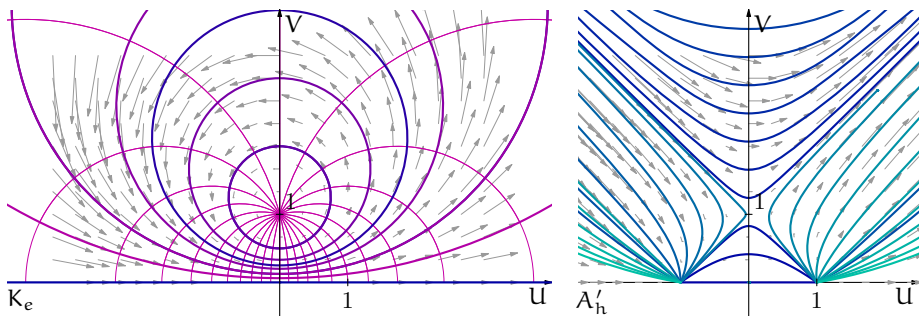


Figure: Elliptic and hyperbolic fix groups of the imaginary units.
In the hyperbolic case there are fixed geometric sets: $\{-1, 1\}$, $(-1, 1)$, \mathbb{R} .

Hyperbolic Case in the Half Space

Cauchy Type Integral

In the **hyperbolic** case we consider *principal* series UIR ρ_σ of $SL_2(\mathbb{R})$. The corresponding decomposition for hyperbolic subgroup A :

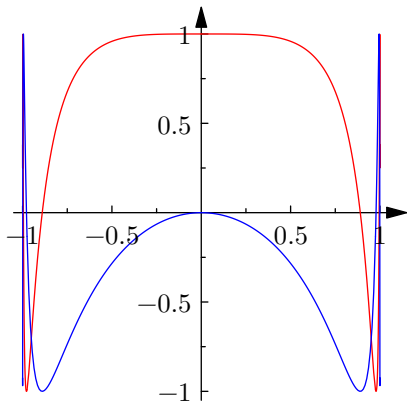
$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = |a| \frac{1}{\sqrt{v}} \begin{pmatrix} v & u \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \cosh t & \sinh t \\ \sinh t & \cosh t \end{pmatrix}$$

For the representation ρ_σ an λ -eigenfunction of the vector field A is:

$$f(x) = \frac{1}{(x^2 - 1)^{\sigma/2}} \left(\frac{x+1}{x-1} \right)^{\lambda/2}.$$

For the value $\lambda = i\sigma$ we have the function:

$$f_0(x) = \frac{1}{(x - \epsilon)^\sigma}.$$



Hyperbolic Wavelets from Double Numbers

The choice of the \mathcal{A} -eigenvector as mother wavelet:

- $f_0 = \delta(x \pm 1)$ —Dirichlet condition.
- $f_0 = \frac{1}{(x - \epsilon)^\sigma} = \left(\frac{x + \epsilon}{x^2 - 1} \right)^\sigma$ —Neumann condition.
- $f_0 = \frac{\chi(1 - x^2)}{(x - \epsilon)^\sigma}$ —space-like and time-like separation, Fig. 2.
- ... (combination of above)

Then we follow the general scheme both for wavelets with complex and double valued wavelets:

- *wavelets* or *coherent states* $v_\sigma(g, z) = \rho_\sigma(g)v_0(z)$.
- d’Alambert integral from the universal *wavelet transforms*

$$\mathcal{W}_\sigma : f(z) \mapsto \mathcal{W}_\sigma f(u) = \langle f(z), \rho_\sigma v_0(u, z) \rangle$$

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Cauchy Type Integral

In the parabolic case we consider UIR ρ_π of $SL_2(\mathbb{R})$. The corresponding decomposition for hyperbolic subgroup N' :

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = |a| \frac{1}{\sqrt{v}} \begin{pmatrix} v & u \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ t & 1 \end{pmatrix}$$

The choice of the N' -eigenvector as mother wavelet:

- $f_0 = \delta(x)$ —previous dependence from the single point.
- $f_0 = \frac{1}{(x - \varepsilon)^\pi} = \frac{1}{x^\pi} + \frac{\varepsilon}{x^{\pi+1}}$ —full Cauchy formula.
- Wavelet generated by the fundamental solution of the heat equation.

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Cauchy-Riemann Equation

from Invariant Fields

Let ρ be a unitary representation of Lie group G with the derived representation $d\rho$ of \mathfrak{g} . Let a mother wavelet w_0 be a null-solution, i.e. $Aw_0 = 0$, for the operator $A = \sum_j a_j d\rho^{X_j}$, where $X_j \in \mathfrak{g}$. Then the wavelet transform $F(g) = \mathcal{W}f(g) = \langle f, \rho(g)w_0 \rangle$ for any f satisfies to:

$$DF(g) = 0, \quad \text{where } D = \sum_j a_j \mathcal{L}^{X_j}.$$

Here \mathcal{L}^{X_j} are left the invariant fields (Lie derivatives) on G corresponding to X_j .

If \mathcal{L}^{X_j} is derived representation of Lie derivative A, N, K (without the matching subgroup) then C-R operator and Laplacian are given by:

$$D = \iota \mathcal{L}^A + \mathcal{L}^X, \quad \text{and} \quad \Delta = D\bar{D} = -\sigma \mathcal{L}^A{}^2 + \mathcal{L}^X{}^2, \quad (9)$$

where X is in the orthogonal complement (with respect to the Killing form) of the corresponding subgroup K, N, A .

Taylor Expansion over Eigenfunctions

Wavelet are decomposable $v_g(x) = \sum_{\alpha} c_{\alpha}(x)\phi_{\alpha}(g)$ over the *complete set* of its eigenfunctions $\phi_{\alpha}(u)$ of the principal subgroup. The C-R operators kill half of them (“negative powers”), only the other half is really needed for the decomposition. Then from the wavelet transform:

$$\langle f(z), v_g(z) \rangle = \left\langle f(z), \sum_{\alpha} c_{\alpha}(z)\phi_{\alpha}(g) \right\rangle = \sum_{\alpha} \phi_{\alpha}(g) \langle f(z), c_{\alpha}(z) \rangle$$

In the **elliptic** case eigenvectors of \mathbf{K} are z^m , $m = 0, 1, 2, \dots$ and the decomposition is the *Taylor series*:

$$f(z) = \sum_0^{\infty} c_n z^n.$$

In the **hyperbolic** case eigenvectors of \mathbf{A} are z^p , $p \in \mathbb{R}_+$ and a Taylor type expansion is given by the integral (not series!)

$$f(z) = \int_0^{\infty} c(p)z^p dp.$$

$SL_2(\mathbb{R})$ Actions on Algebras/Moduli

Let $\mathfrak{a} \in \mathfrak{A}$ with $\mathfrak{sp} \mathfrak{a} \in \bar{\mathbb{D}}$ be fixed in a Banach algebra \mathfrak{A} with the unit e , then

$$g : \mathfrak{a} \mapsto g \cdot \mathfrak{a} = (\bar{\alpha}\mathfrak{a} - \bar{\beta}e)(\alpha e - \beta\mathfrak{a})^{-1}, \quad g \in SL_2(\mathbb{R}) \quad (10)$$

is a well defined $SL_2(\mathbb{R})$ action on a subset $\mathbb{A} = \{g \cdot \mathfrak{a} \mid g \in SL_2(\mathbb{R})\} \in \mathfrak{A}$, i.e. \mathbb{A} is a $SL_2(\mathbb{R})$ -homogeneous space.

Define *resolvent* function $R(g, \mathfrak{a}) : \mathbb{A} \rightarrow \mathfrak{A}$:

$$R(g, \mathfrak{a}) = (\alpha e - \beta\mathfrak{a})^{-1} \quad \text{then} \quad R_1(g_1, \mathfrak{a})R_1(g_2, g_1^{-1}\mathfrak{a}) = R_1(g_1g_2, \mathfrak{a}). \quad (11)$$

Thus we can linearise (10) in $C(\mathbb{A}, M)$, for a **left \mathfrak{A} -module** M (e.g. $M = \mathfrak{A}$):

$$\begin{aligned} \rho_{\mathfrak{a}}(g_1) : f(g^{-1} \cdot \mathfrak{a}) &\mapsto R(g_1^{-1}g^{-1}, \mathfrak{a})f(g_1^{-1}g^{-1} \cdot \mathfrak{a}) \\ &= (\alpha'e - \beta'\mathfrak{a})^{-1} f\left(\frac{\bar{\alpha}' \cdot \mathfrak{a} - \bar{\beta}'e}{\alpha'e - \beta'\mathfrak{a}}\right). \end{aligned}$$

Wavelet Transform in Module Spaces

For any $x \in M'$ define a constant M' -valued *vacuum vector*

$$v_x(a) = x \otimes v_0(g) \in C(\mathbb{A}, M)$$

The *wavelet transform* associated with v_x is defined by the same formula:

$$\mathcal{W}_m f(g) = \langle f, \rho_a(g) v_x \rangle$$

which is vector versions of Cauchy or Bergman integral. It maps $L_2(\mathbb{A})$ to $C(SL_2(\mathbb{R}), \mathbb{C})$.

An integral representation for a covariant calculus Φ is provided by the inverse wavelet transform:

$$\mathcal{M} : f(g) \mapsto f(a, v_0) = \int_G f(g) \rho_a(g) dg v_0$$

The *Riesz-Dunford calculus* is given by

$$\Phi : f \mapsto \mathcal{W}_1 f(a) = \frac{1}{2\pi i} \int_{\mathbb{T}} f(z) \frac{1}{a - z} dz$$

where $M = \mathfrak{A}$ and $v_0 = e$.

Jet Spaces

Definition

Two holomorphic functions have n th order contact in a point if their value and their first n derivatives agree at that point.

A point $(z, u^{(n)}) = (z, u, u_1, \dots, u_n)$ of the jet space $\mathbb{J}^n \sim \mathbb{D} \times \mathbb{C}^n$ is the equivalence class of holomorphic functions having n th contact at the point z .

For a fixed n each holomorphic function $f : \mathbb{D} \rightarrow \mathbb{C}$ has n th prolongation (or n -jet) $jf : \mathbb{D} \rightarrow \mathbb{C}^{n+1}$ defined as follows:

$$jf(z) = (f(z), f'(z), \dots, f^{(n)}(z)).$$

Jet spaces and prolongations were introduced by S. Lie to study symmetries of the differential equations.

Prolongation of Representations

The representation ρ_m of the group $SL_2(\mathbb{R})$ in $B_m(\mathbb{D})$ could be prolonged to a representation $\rho_m^{(n)}$ of $SL_2(\mathbb{R})$ by a transformation $\rho_1^{(n)} : jf(z) \mapsto \rho_1^{(n)} jf(z)$ of the jet space \mathbb{J}^n :

$$\rho_m^{(n)}(g) : (z, u, \dots, u_n) \mapsto (z(g), u(g), \dots, u_n(g)), \quad \text{where } z(g) = \frac{\bar{\alpha}z - \bar{\beta}}{-\beta z + \alpha},$$

and $u_k(g)$ is the k th derivative of $\rho_m u$ at the point $z(g)$.

From the definition: j intertwines ρ_1 and $\rho_1^{(n)}$:

$$j\rho_1(g) = \rho_1^{(n)}(g)j \quad \text{for all } g \in SL_2(\mathbb{R}).$$

Proposition

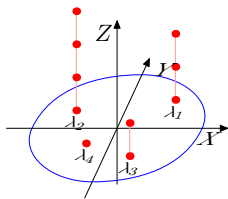
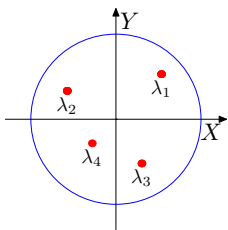
Let α is a Jordan block of a length k for $\lambda = 0$, and x be its root vector of order k i.e. $\alpha^{k-1}x \neq \alpha^k x = 0$. Then $\rho_{\alpha, m}$ on v_x is equivalent to ρ_m^k .

Spectrum of a Jordan block

Because of the transitive group of inner automorphisms, which could send any $\lambda \in \mathbb{D}$ to 0 , we got the complete characterisation of $\rho_{\mathbf{a}}$ for matrices.

Proposition (Jordan normal form)

Representation $\rho_{\mathbf{a}}$ is equivalent to a direct sum of the prolongations $\rho_{\mathbf{m}}^{(k)}$ of $\rho_{\mathbf{m}}$ in the k th jet space \mathbb{J}^k intertwined with inner automorphisms. Consequently the spectrum of \mathbf{a} (defined via the functional calculus $\Phi = \mathcal{W}_{\mathbf{m}}$) consists of exactly n pairs (λ_i, k_i) , $\lambda_i \in \mathbb{D}$, $k_i \in \mathbb{Z}_+$, $1 \leq i \leq n$.



Traditional (left) and covariant (right) spectra of the matrix:
 $\mathbf{a} = J_3(\lambda_1) \oplus J_4(\lambda_2) \oplus J_2(\lambda_3) \oplus J_1(\lambda_4)$.

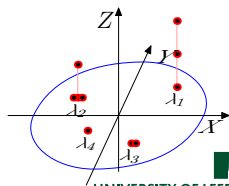
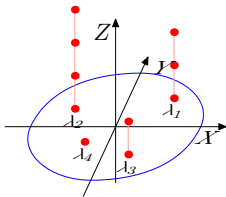
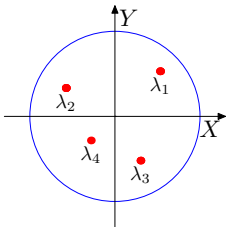
Theorem (Spectral mapping)

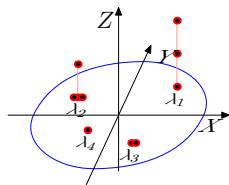
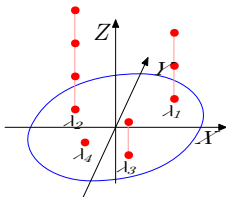
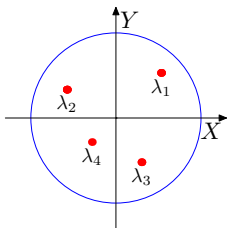
Let $\phi : \mathbb{D} \rightarrow \mathbb{D}$ be a holomorphic map, let us define $[\phi_* f](z) = f(\phi(z))$ and its prolongation $\phi_*^{(n)}$ onto the jet space \mathbb{J}^n . Its associated action $\rho_1^k \phi_*^{(n)} = \phi_*^{(n)} \rho_1^n$ on the pairs (λ, k) is given by the formula:

$$\phi_*^{(n)}(\lambda, k) = \left(\phi(\lambda), \left[\frac{k}{\deg_\lambda \phi} \right] \right),$$

where $\deg_\lambda \phi$ denotes the degree of zero of the function $\phi(z) - \phi(\lambda)$ at the point $z = \lambda$ and $[x]$ denotes the integer part of x . Then

$\text{sp } \phi(a) = \phi_*^{(n)} \text{sp } a$ (which is actually known for Jordan blocks).





Two first pictures illustrate the traditional and new spectra of the matrix:

$$\mathbf{a} = J_3(\lambda_1) \oplus J_4(\lambda_2) \oplus J_2(\lambda_3) \oplus J_1(\lambda_4).$$

The traditional spectrum is *the same* for many essentially different (even by dimensionality) matrices, e.g. for \mathbf{a} above and

$$\mathbf{a}_1 = \text{diag}(\lambda_1, \lambda_2, \lambda_3, \lambda_4).$$

The third picture shows *spectral mapping* for a function ϕ which fixes all four eigenvalues $\lambda_1, \dots, \lambda_4$ and has such *orders of zeros*: the order 1 at λ_1 , exactly the order 3 at λ_2 , an order at least 2 at λ_3 , and finally any order at λ_4 .

Spectrum of a Perturbation of a Jordan Block

Theorem (V.B. Lidskii, 1965)

Eigenvalues of the perturbed matrix $J_n + c^n K$ admit the expansion

$$\lambda_j = c\xi^{1/n} + o(c),$$

where $\xi^{1/n}$ represents all n -th roots of certain $\xi \in \mathbb{C}$.

Theorem (E.B. Davies, M. Hager, 2006)

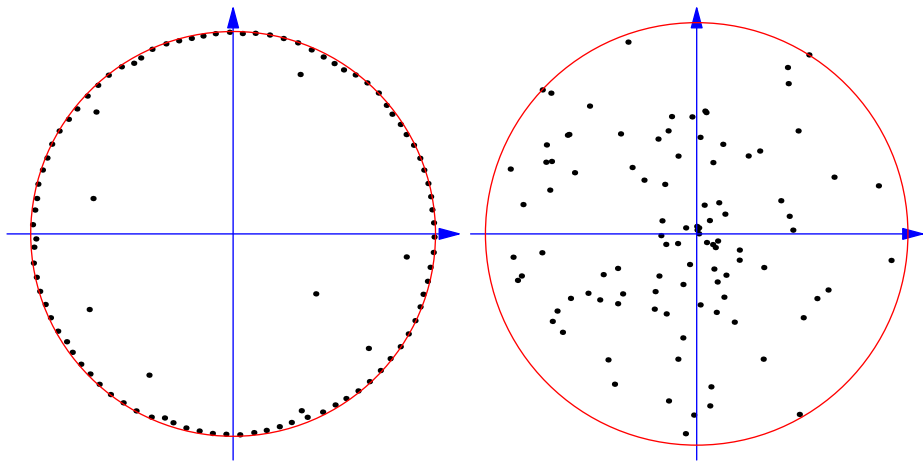
Let $M = J + c^n K$ where J is the standard $n \times n$ Jordan matrix, $0 < c < 1$ and K is a random matrix with independent Gaussian entries.

Then for any $\varepsilon > 0$ with probability that converges to 1 as $n \rightarrow \infty$, the proportion of the eigenvalues that lie in any annulus

$$\{z : c - \varepsilon < |z| < c + \varepsilon\}$$

converges to 1. The remaining eigenvalues lie inside the annulus.

Is It “Spectral Instability”?



(a) (b)

Figure: (a) The spectrum of the perturbation $J_{100} + C$ of the Jordan block, see a movie and another one.

(b) The spectrum of the random matrix C

Spectral Distance

The following notion links functional calculus with functional models:

Definition

The spectral distance between two points in the jet space is equal to distance in $H_2(\mathbb{T})$ between two Blaschke products with respective zeroes.

Theorem

Let $n = 2$ and let $\lambda_1(t)$ and $\lambda_2(t)$ are eigenvalues of the matrix $J_2 + c^2 \cdot K$ with some matrix K . Then

$$|\lambda_1(c)| + |\lambda_2(c)| = O(c), \text{ however } |\lambda_1(c) + \lambda_2(c)| = O(c^2). \quad (12)$$

The spectral distance from the 1-jet at 0 to two 0-jets at points λ_1 and λ_2 bounded only by the first condition in (12) is $O(c^2)$.

However the spectral distance between J_2 and $J_2 + c^2 \cdot K$ is $O(c^4)$.

Hyperbolic Functional Calculus

Elliptic case

Variable $z = e^{it} = \cos t + i \sin t$

Taylor series $f(z) = \sum_0^\infty c_n z^n$

Multiplication by e^{it}

Shift S_d on \mathbb{Z}

Root vector $v_k: S_d^k v_k = 0$

Calculus $f(S_d)v = \sum_0^{k-1} c_n S_d^n v_k$

$H_n = \text{Lin}\{S_d^n v_k, n = 1, \dots, k-1\}$

Jordan block on H_n

Hyperbolic case

Variable $z = e^{\epsilon t} = \cosh t + \epsilon \sinh t$

Taylor integral $f(z) = \int_0^\infty c(p) z^p dp$

Multiplication by $e^{\epsilon t}$

Shift S_c on \mathbb{R}

Root vector $v_a: S_c^a v_a = 0$

Calculus $f(S_c) = \int_0^a c(t) S_c^t dt v_a$

$H_a = \overline{\text{Lin}}\{S_c^t v_a, t \in [0, a]\}$

Integral operator on H_a

Polynomially Bounded Operators

Standard for \mathfrak{a} with $\text{sp } \mathfrak{a} \in \bar{\mathbb{D}}$ and $\|\mathfrak{a}^k\| < Ck^p$ to consider *power bounded* $r\mathfrak{a}$, where $0 < r < 1$, and its H_∞ calculus. Further properties of an operator are recovered from consideration of the limit $r \rightarrow 1$.

A *better regularisation*, $\mathfrak{a}^k \rightarrow \mathfrak{a}^k/k^p$ rather than $\mathfrak{a}^k \rightarrow r^k \mathfrak{a}^k$, is achieved in the covariant calculus framework (although algebra homomorphism is *completely* lost).

Since norm of $f(z) = \sum_0^\infty c_k z^k$ in B_m is equivalent to $\sum_0^\infty c_k^2/k^{m-1}$ for polynomially bounded \mathfrak{a} the resolvent $R(z, \mathfrak{a})$ belongs to any B_m with $m > 2(p+1)$.

Define a representation of $SL_2(\mathbb{R})$ in $B_m(\mathbb{D} \times \mathbb{A}, M)$ by:

$$\rho_m' : f(\mathfrak{u}, \mathfrak{a}) \mapsto \frac{1}{(\bar{\beta}\mathfrak{u} + \alpha)^{m-1}(\alpha\mathfrak{e} - \beta\mathfrak{a})} f\left(\mathfrak{u}, \frac{\bar{\alpha}\mathfrak{a} - \beta\mathfrak{e}}{\alpha\mathfrak{e} - \beta\mathfrak{a}}\right).$$

It is generated by the discrete series representation of $SL_2(\mathbb{R})$ with the *lowest weight* m .

Calculus of Polynomially Bounded Operators in Bergman Spaces

For the vacuum vector $v_0(u, a) \equiv x$ in $B_m(\mathbb{D} \times \mathbb{A}, M)$, where $(x \in M)$, the corresponding functional calculus is given by the integral:

$$f(g \cdot a) = \int_{\mathbb{D}} \frac{f(u)}{(\beta \bar{u} + \bar{\alpha})^{m-1} (\bar{\alpha} e - \bar{\beta} a)} \frac{du}{(1 - |u|^2)^{m-2}}.$$

For Jordan k -blocks with $|\lambda_i| = 1$ it is equivalent to k -prolongation of ρ_m' . Since $B_m \subset B_n$ if $m \leq n$ the existence ρ_m -calculus implies ρ_n calculus. The minimal m such that B_m covariant calculus exists measures the order of operator.

Several Variables Spectral Theory

For a *joint spectrum* of n -tuple of operators we have many alternatives:

- Weyl functional calculus through the Heisenberg group \mathbb{H}^n acting in $L_2(\mathbb{R}^n)$;
- Segal-Bargmann type functional calculus through the Heisenberg group \mathbb{H}^n acting in $L_2(\mathbb{C}^n)$;
- several complex variables through groups of automorphisms of unit ball or polydisk in \mathbb{C}^n . However this is suitable mainly for commuting n -tuples;

Or the Clifford analysis through the Möbius group of conformal maps of \mathbb{C}^n . The Clifford algebra $\mathcal{C}l(n)$ is spanned by $1, e_1, e_2, \dots, e_n$ with relations

$$e_k^2 = -1 \quad \text{and} \quad e_k e_j = -e_j e_k \quad \text{for } k \neq j.$$

Similarly to complex analysis we could derive a *Cauchy kernel* in $B(\mathbb{H}) \otimes \mathcal{C}l(n)$ (cf. resolvent):

$$R(A_1, A_2, \dots, A_n; \lambda_1, \lambda_2, \dots, \lambda_n) = \left(\sum_{k=1}^n e_k A_k - \sum_{k=1}^n e_k \lambda_k I \right)^{-1}$$

Joint Spectra of Pauli Matrices

Example

Let $J_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ and $J_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ be the Pauli matrices.

The Cauchy kernel is:

$$\frac{-\lambda_1^2 + \lambda_2^2 + 2\lambda_1\lambda_2 e_1 e_2}{(\lambda_1^2 + \lambda_2^2)^2} \begin{pmatrix} (-1 - \lambda_1)e_1 - \lambda_2 e_2 & e_2 \\ e_2 & (1 - \lambda_1)e_1 - \lambda_2 e_2 \end{pmatrix}$$

Comparison of different *joint* spectra for J_1 and J_2 :

Clifford spectrum (invertibility of $A - \lambda I$)	$\mathbf{sp}_C(J_1, J_2) = \{(0, 0)\},$
Weyl spectrum (support of operator-valued distribution)	$\mathbf{sp}_W(J_1, J_2) = \mathbb{D},$
Möbius spectrum (support of the intertwining operator)	$\mathbf{sp}_M(J_1, J_2) = \{\rho_1, \rho_1^{(1)}\}.$

